

Cosmology based on Finsler geometry

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Outline

- 1 Motivation
- 2 Finsler cosmology
- 3 Tensorial Finsler cosmologies
- 4 Conclusion

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Why cosmology based on Finsler geometry?

- Modify spacetime geometry to address open problems:
 - Origin of dark matter and dark energy.
 - Homogeneity of the cosmic microwave background and inflation.
 - Fly-by anomaly in the solar system.

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 - Divide tangent spaces into space-, time-, lightlike vectors.
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 - Distinguish curves corresponding to physical trajectories.
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 - Determine trajectories of freely falling test masses.
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 - **Gravity theory on Finsler spacetimes exists.**

Finsler gravity action and field equations

- Spacetime manifold M with Finsler function $F : TM \rightarrow \mathbb{R}^+$.

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- Finsler gravity action:

$$S_G = \int_{\Sigma} \mathcal{R} \text{Vol}(G|_{\Sigma}).$$

- Σ : Unit tangent bundle $TM|_{F=1}$.
- G : Sasaki metric on TM .
- \mathcal{R} : Scalar curvature of Cartan non-linear connection.

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- Gravitational field equations by variation with respect to F :

$$-\frac{1}{F^2} \left\{ 6\mathcal{R} + G^{ab} \left[\nabla_a^v \nabla_b^v \mathcal{R} + 2F^2 J^c_a \nabla_b^h \mathcal{S}_c + 2\nabla_a^v (\mathbf{S}^c \nabla_c^h \mathcal{S}_b) \right] \right\} = \mathcal{T}.$$

- a, b, c, \dots : Coordinate indices $0, \dots, 7$ on TM .
- J^a_b : Tangent structure.
- \mathbf{S}^a : Geodesic spray.
- \mathcal{S}_a : Landsberg covector.
- ∇^h, ∇^v : Horizontal and vertical Berwald derivative.
- \mathcal{T} : Energy-momentum scalar.

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Finsler length function

- Finsler function $F : TM \rightarrow \mathbb{R}^+$.
- Finsler length functional for $\gamma : \mathbb{R} \rightarrow M$:

$$\ell_{t_1}^{t_2}[\gamma] = \int_{t_1}^{t_2} F(\gamma(t), \dot{\gamma}(t)) dt .$$

- Parametrization invariance requires homogeneity:

$$F(x, \lambda y) = \lambda F(x, y) \quad \forall \lambda > 0 .$$

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Connection and curvature

- Split $TTM = HTM \oplus VTM$ of the double tangent bundle.
- Projectors \mathbf{h} , \mathbf{v} onto subbundles.
- Curvature tensor $R = -N_{\mathbf{h}}$ defined via Nijenhuis tensor.
- Curvature scalar defined from curvature tensor.

Cosmological coordinates on TM [MH '15]

- Spherical coordinates t, r, ϑ, φ on M .
- Coordinates y, u, v, w on each $T_x M$:

$$y\partial_t + w \left[\cos u \sqrt{1 - kr^2} \partial_r + \frac{\sin u}{r} \left(\cos v \partial_\vartheta + \frac{\sin v}{\sin \vartheta} \partial_\varphi \right) \right] \in T_x M.$$

Cosmological symmetry

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Cosmologically symmetric Finsler spacetime

- Symmetry under rotations and translations (six vector fields).
- Most general Finsler function: $F(t, y, w)$.
- Homogeneity condition: $F(t, \lambda y, \lambda w) = \lambda F(t, y, w)$.
- Express Finsler function as $F(t, y, w) = y \tilde{F}(t, w/y)$.

Observer trajectories

- Tangent vectors are future unit timelike vectors: $F = 1$.
- Future unit timelike vectors form shell in each $T_x M$.
- Introduce suitable coordinates on these shells.

Observer space coordinates

Observer trajectories

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- Introduce suitable coordinates on these shells.

Observer space coordinates [MH '15]

- Introduce coordinates:

$$T = t, R = r, \Theta = \vartheta, \Phi = \varphi, Y = y\tilde{F}\left(t, \frac{w}{y}\right), U = u, V = v, W = w/y.$$

⇒ Unit tangent bundle has $Y = 1$.

⇒ Light cone has $Y = 0$.

Geodesics on cosmological background

- $\gamma : \mathbb{R} \rightarrow M$ minimizes Finsler length functional.
- $\Leftrightarrow \gamma$ satisfies second order ODE.
- $\Leftrightarrow \dot{\gamma} : \mathbb{R} \rightarrow TM$ satisfies first order ODE.
- $\Leftrightarrow \dot{\gamma}$ is integral curve of vector field \mathbf{S} called **geodesic spray**:

$$\mathbf{S} = \frac{Y}{\tilde{F}} \left(\partial_T + W \cos U \sqrt{1 - kR^2} \partial_R + \frac{W \sin U \cos V}{R} \partial_\Theta + \frac{W \sin U \sin V}{R \sin \Theta} \partial_\Phi - \frac{W \sin U \sqrt{1 - kR^2}}{R} \partial_U - \frac{W \sin U \sin V}{R \tan \Theta} \partial_V - \frac{\partial_T \partial_W \tilde{F}}{\partial_W \partial_W \tilde{F}} \partial_W \right).$$

- Coordinate Y is constant along Finsler geodesics.

Geodesics on cosmological background

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- $\Leftrightarrow \dot{\gamma} : \mathbb{R} \rightarrow TM$ satisfies first order ODE.
- $\Leftrightarrow \dot{\gamma}$ is integral curve of vector field \mathbf{S} called geodesic spray.
- Coordinate Y is constant along Finsler geodesics.
- **Radial geodesic** given by $U = 0$:

$$\mathbf{S}|_{U=0} = \frac{Y}{\tilde{F}} \left(\partial_T + W\sqrt{1 - kR^2}\partial_R - \frac{\partial_T\partial_W\tilde{F}}{\partial_W\partial_W\tilde{F}}\partial_W \right).$$

- **Co-moving geodesic** given by $W = 0$:

$$\mathbf{S}|_{W=0} = \frac{Y}{\tilde{F}} \left(\partial_T - \frac{\partial_T\partial_W\tilde{F}}{\partial_W\partial_W\tilde{F}}\partial_W \right).$$

Kinetic theory of fluids [Ehlers '71], [Sarbach, Zannias '13]

- Consider fluid as constituted by point particles.
- Particles follow piecewise geodesics between collisions.
- Continuum limit described by density $\phi : TM|_{Y=1} \rightarrow \mathbb{R}^+$.
- Collisionless fluid satisfies Liouville equation $\mathcal{L}_S \phi = 0$.

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Cosmologically symmetric Finsler fluids [MH '15]

- Most general cosmologically symmetric fluid: $\phi = \phi(T, W)$.
- Liouville equation: $\partial_T \phi \partial_W \partial_W \check{F} = \partial_W \phi \partial_T \partial_W \check{F}$.

Finsler gravity [Pfeifer, Wohlfarth '11]

- Action:

$$S_G = \int_{\Sigma} \mathcal{R} \text{Vol}(G|_{\Sigma}).$$

- Field equations:

$$-\frac{1}{F^2} \left\{ 6\mathcal{R} + G^{ab} \left[\nabla_a^v \nabla_b^v \mathcal{R} + 2F^2 J^c_a \nabla_b^h \mathcal{S}_c + 2\nabla_a^v \left(\mathbf{S}^c \nabla_c^h \mathcal{S}_b \right) \right] \right\} = \mathcal{T}.$$

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Cosmological dynamics

- Structure of cosmological equations: $\mathcal{G}[\tilde{F}](T, W) = \mathcal{T}[\tilde{F}, \phi](T, W)$.
- Difficulties:
 - Geometry scalar \mathcal{G} is complicated even for cosmology.
 - No “standard construction” for \mathcal{T} of non-metric kinetic fluid.

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Geometry

- Tensor field: metric $g_{\mu\nu}$.
- Cosmology: FLRW metric $g = -dt \otimes dt + a^2(t)\gamma_{ij}[\kappa]dx^i \otimes dx^j$.
- Finsler function:

$$F(x, y) = \sqrt{|g_{\mu\nu}y^\mu y^\nu|} = Y\sqrt{|1 - a^2(T)W^2|}$$

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Gravitational dynamics

- Geometry scalar:

$$\mathcal{G} = \frac{6}{a^2(1 - W^2 a^2)} \left(a\ddot{a} - 2\dot{a}^2 - 2\kappa + W^2 a^3 \ddot{a} \right).$$

⇒ Reproduce structure of Friedmann equations.

Length measure with one-forms

Ingredients

- Tensor fields: metric $g_{\mu\nu}$, one-form A_μ .
- Cosmology:
 - FLRW metric $g = -dt \otimes dt + a^2(t)\gamma_{ij}[\kappa]dx^i \otimes dx^j$.
 - Hypersurface normal $A = b(t)dt$.

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Randers length measure [Randers '41]

$$F(x, y) = \sqrt{|g_{\mu\nu}y^\mu y^\nu|} + A_\mu y^\mu = Y\sqrt{|1 - a^2(T)W^2|} + Yb(T)$$

Bogoslovsky length measure [Bogoslovsky '77]

$$F(x, y) = (A_\mu y^\mu)^q \left(\sqrt{|g_{\mu\nu}y^\mu y^\nu|} \right)^{1-q} = Yb^q(T) \left(\sqrt{|1 - a^2(T)W^2|} \right)^{1-q}$$

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Summary

- Finsler spacetimes:
 - Based on Finsler length function.
 - Make use of tensors on the tangent bundle.
 - Generalize standard notions of causality, observers and gravity.
- Finsler cosmology:
 - Geometry defined by function $\tilde{F}(T, W)$.
 - Simple form of geodesic equation.
 - Simple equation of motion for fluid dynamics.
 - Gravitational field equations are rather complicated.
 - Simplified models can be derived from tensorial geometries.

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Outlook

- Construct energy-momentum scalar for kinetic Finsler fluid.
- Find non-metric solutions for Finsler cosmology.
- Calculate cosmological parameters from Finsler geometry.

Conclusion

- Kinetic theory on the tangent bundle:
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