

Gravitational waves in teleparallel theories of gravity

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(with Martin Krššák, Christian Pfeifer, Jackson Levi Said, Ulbossyn Ualikhanova)

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Center of Excellence "The Dark Side of the Universe"



The third Zeldovich meeting - 25. April 2018



THE THIRD ZELDOVICH MEETING SCIENTIFIC OBJECTIVES



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Visitors: 108878748
We have 3 guests online

International Center for Relativistic Astrophysics Network (ICRANet) together with the National Academy of Sciences of Belarus organize an international conference to be held in Minsk, Belarus in April 23-27, 2018. Participation from neighboring countries [such as Estonia](#), Latvia, Lithuania, Poland, Russia and Ukraine as well as from Balkan countries, Eastern and Western Europe and the Americas is expected. Exceptionally wide research interests of Ya. B. Zeldovich ranging from chemical physics, elementary particle and nuclear physics to astrophysics and cosmology provide the topics to be covered at the conference:

- Early cosmology, large scale structure, cosmic microwave background;
- Neutron stars, black holes, gamma-ray bursts, supernovae, hypernovae;
- Ultra high energy particles;
- Gravitational waves.

Many lecturers at the conference will be the members of the world-famous scientific school in astrophysics and cosmology, founded by Ya. B. Zeldovich, who now became leading scientists in these fields in many countries worldwide including Germany, Italy, USA and Russia.

This conference will follow a very successful international conferences in honor of Ya. B. Zeldovich, held in Minsk in 2009 and in 2014.

For suggestions&comments write to the [Webmaster](#)

Zeldovich in Estonia (end of 1970s)



Summer school on cosmology, Tõravere 1962



- 1 Introduction
- 2 Waves in torsion gravity
- 3 Waves in non-metricity gravity
- 4 Conclusion

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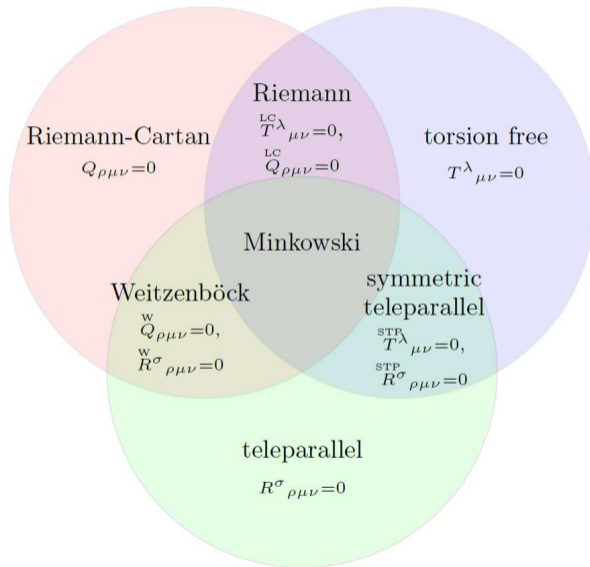
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 - Accelerating phases in the history of the Universe?
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 - Describes gravity as gauge theory of the translation group.
 - First order action, second order field equations.
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- Gravity formulated as gauge theories.

Overview of geometries



- Complex double null basis of the tangent bundle:

$$l = \partial_t + \partial_z, \quad n = \frac{\partial_t - \partial_z}{2}, \quad m = \frac{\partial_x + i\partial_y}{\sqrt{2}}, \quad \bar{m} = \frac{\partial_x - i\partial_y}{\sqrt{2}}.$$

Newman-Penrose formalism

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- Consider plane null wave with $k_\mu = -\omega l_\mu$ and $u = t - z$:

$$h_{\mu\nu} = H_{\mu\nu} e^{ik_\mu x^\mu} = H_{\mu\nu} e^{i\omega u}.$$

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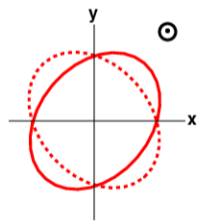
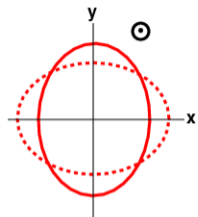
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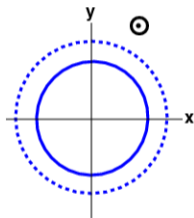
- Riemann tensor determined by “electric” components:

$$\begin{aligned} \Psi_2 &= -\frac{1}{6} R_{nlnl} = \frac{1}{12} \ddot{h}_{ll}, & \Psi_3 &= -\frac{1}{2} R_{nl n\bar{m}} = \frac{1}{4} \ddot{h}_{l\bar{m}}, \\ \Psi_4 &= -R_{n\bar{m}n\bar{m}} = \frac{1}{2} \ddot{h}_{\bar{m}\bar{m}}, & \Phi_{22} &= -R_{nmn\bar{m}} = \frac{1}{2} \ddot{h}_{m\bar{m}}. \end{aligned}$$

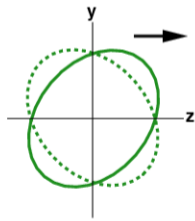
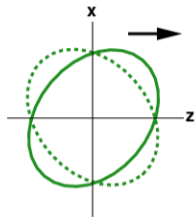
Polarisations of gravitational waves



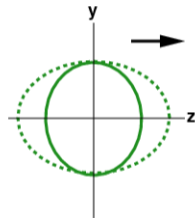
$\Psi_4, \bar{\Psi}_4$



Φ_{22}



$\Psi_3, \bar{\Psi}_3$



Ψ_2

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- Derived quantities:
 - Frame field $e_a = e_a{}^\mu \partial_\mu$ with $\iota_{e_a} \theta^b = \delta_a^b$.
 - Metric $g_{\mu\nu} = \eta_{ab} \theta^a{}_\mu \theta^b{}_\nu$.
 - Volume form $\theta d^4x = \theta^0 \wedge \theta^1 \wedge \theta^2 \wedge \theta^3$.
 - Levi-Civita connection

$$\overset{\circ}{\omega}{}_{ab} = -\frac{1}{2} (\iota_{e_b} \iota_{e_c} d\theta_a + \iota_{e_c} \iota_{e_a} d\theta_b - \iota_{e_a} \iota_{e_b} d\theta_c) \theta^c.$$

- Torsion $T^a = d\theta^a + \overset{\circ}{\omega}{}^a{}_b \wedge \theta^b$.

Field content and geometry

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- Gauge fixing

- Perform local Lorentz transformation:

$$\theta'^a = \Lambda^a{}_b \theta^b, \quad \dot{\omega}'^a{}_b = \Lambda^a{}_c \dot{\omega}^c{}_d \Lambda_b{}^d + \Lambda^a{}_c d\Lambda_b{}^c.$$

⇒ Weitzenböck gauge: set $\dot{\omega}^a{}_b \equiv 0$.

Most general action and corresponding field equations

- Most general action:

$$S = \frac{1}{2\kappa^2} \int d^4x e (c_1 T^{\mu\nu\rho} T_{\mu\nu\rho} + c_2 T^{\mu\nu\rho} T_{\rho\nu\mu} + c_3 T^\mu{}_{\mu\rho} T_\nu{}^{\nu\rho}) .$$

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- Field tensors:

- Symmetric perturbation part:

$$F^{\mu\rho\sigma} = (2c_1 + c_2) (\partial^\sigma \phi^{\mu\rho} - \partial^\rho \phi^{\mu\sigma}) + c_3 [(\partial^\sigma \phi^\alpha{}_\alpha - \partial_\alpha \phi^{\alpha\sigma}) \eta^{\mu\rho} - (\partial^\rho \phi^\alpha{}_\alpha - \partial_\alpha \phi^{\alpha\rho}) \eta^{\mu\sigma}] .$$

- Antisymmetric perturbation part:

$$B^{\mu\rho\sigma} = (2c_1 - c_2) (\partial^\sigma a^{\mu\rho} - \partial^\rho a^{\mu\sigma}) + (2c_2 + c_3) \partial^\mu a^{\sigma\rho} .$$

Newman-Penrose decomposition

Field equations expressed in Newman-Penrose basis

$$0 = E_{nn} = (2c_1 + c_2 + c_3)\partial_n^2\phi_{nl} + 2c_3\phi_{m\bar{m}} + (2c_1 + c_2 + c_3)\partial_n^2 a_{nl},$$

$$0 = E_{mn} = (2c_1 + c_2)\partial_n^2\phi_{ml} + (2c_1 - c_2)\partial_n^2 a_{ml},$$

$$0 = E_{\bar{m}n} = (2c_1 + c_2)\partial_n^2\phi_{\bar{m}l} + (2c_1 - c_2)\partial_n^2 a_{\bar{m}l},$$

$$0 = E_{nm} = -c_3\partial_n^2\phi_{lm} - (2c_2 + c_3)\partial_n^2 a_{lm},$$

$$0 = E_{n\bar{m}} = -c_3\partial_n^2\phi_{l\bar{m}} - (2c_2 + c_3)\partial_n^2 a_{l\bar{m}}$$

$$0 = E_{m\bar{m}} = -c_3\partial_n^2\phi_{ll},$$

$$0 = E_{ln} = (2c_1 + c_2)\partial_n^2\phi_{ll}, .$$

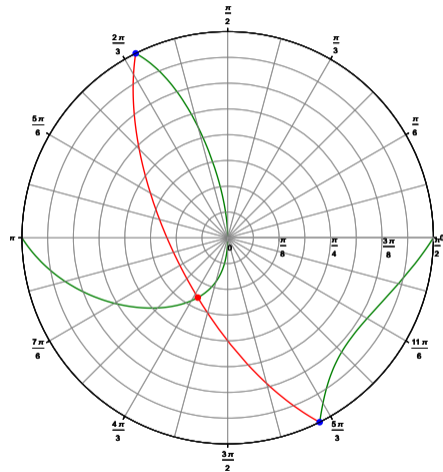
Gravitational wave polarisations

$$c_1 = \sin \theta \cos \phi$$

$$c_2 = \sin \theta \sin \phi$$

$$c_3 = \cos \theta$$

- N_2
- N_3
- III_5
- II_6



Outline

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Field content and geometry

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$$\overset{\circ}{\Gamma}{}^{\rho}{}_{\mu\nu} = \frac{1}{2}g^{\rho\sigma}(\partial_{\mu}g_{\sigma\nu} + \partial_{\nu}g_{\mu\sigma} - \partial_{\sigma}g_{\mu\nu}).$$

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- Gauge fixing

- Perform local coordinate transformation:

$$g'_{\mu\nu} = \frac{\partial x^{\alpha}}{\partial x'^{\mu}} \frac{\partial x^{\beta}}{\partial x'^{\nu}} g_{\alpha\beta}, \quad \overset{\times}{\Gamma}'{}^{\rho}{}_{\mu\nu} = \frac{\partial x^{\alpha}}{\partial x'^{\mu}} \frac{\partial x^{\beta}}{\partial x'^{\nu}} \frac{\partial x'^{\rho}}{\partial x^{\gamma}} \overset{\times}{\Gamma}{}^{\gamma}{}_{\alpha\beta} + \frac{\partial^2 x^{\alpha}}{\partial x'^{\mu} \partial x'^{\nu}} \frac{\partial x'^{\rho}}{\partial x^{\alpha}}.$$

⇒ Coincident gauge: set $\overset{\times}{\Gamma}{}^{\rho}{}_{\mu\nu} \equiv 0 \Rightarrow Q_{\rho\mu\nu} = \partial_{\rho} g_{\mu\nu}$.

Most general action and corresponding field equations

- Most general action:

$$S = - \int d^4x \frac{\sqrt{-g}}{2} \left[c_1 Q^\alpha{}_{\mu\nu} + c_2 Q_{(\mu}{}^\alpha{}_{\nu)} + c_3 Q^\alpha g_{\mu\nu} + c_4 \delta_{(\mu}^\alpha \tilde{Q}_{\nu)} + \frac{c_5}{2} \left(\tilde{Q}^\alpha g_{\mu\nu} + \delta_{(\mu}^\alpha Q_{\nu)} \right) \right]$$

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- Linearized field equations:

$$0 = 2c_1 \eta^{\alpha\sigma} \partial_\alpha \partial_\sigma h_{\mu\nu} + c_2 \eta^{\alpha\sigma} (\partial_\alpha \partial_\mu h_{\sigma\nu} + \partial_\alpha \partial_\nu h_{\sigma\mu}) + 2c_3 \eta_{\mu\nu} \eta^{\tau\omega} \eta^{\alpha\sigma} \partial_\alpha \partial_\sigma h_{\tau\omega} \\ + c_4 \eta^{\omega\sigma} (\partial_\mu \partial_\omega h_{\nu\sigma} + \partial_\nu \partial_\omega h_{\mu\sigma}) + c_5 \eta_{\mu\nu} \eta^{\omega\gamma} \eta^{\alpha\sigma} \partial_\alpha \partial_\omega h_{\sigma\gamma} + c_5 \eta^{\omega\sigma} \partial_\mu \partial_\nu h_{\omega\sigma}.$$

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- Terms involving c_1 and c_3 do not contribute for a null wave $\square h_{\mu\nu} = 0$.

Field equations expressed in Newman-Penrose basis

$$0 = E_{nn} = -2(c_2 \ddot{h}_{ln} + c_4 \ddot{h}_{nl} + c_5 \ddot{h}_{nl} - c_5 \ddot{h}_{m\bar{m}}),$$

$$0 = E_{mn} = E_{nm} = -(c_2 + c_4) \ddot{h}_{lm},$$

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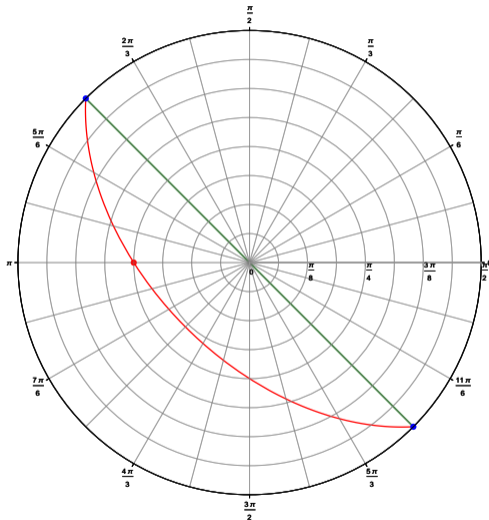
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- Results:
 - Gravitational waves propagate at the speed of light (not shown in this talk).
 - Polarisation classes N_2 , N_3 , III_5 , II_6 : tensor modes always exist, maybe more.

Acknowledgments

- Estonian Research Council:



- European Regional Development Fund:



European Union
European Regional
Development Fund



Investing
in your future

- COST Action CA15117 (CANTATA),
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Acknowledgments

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Teleparallel gravity workshop



June 25-29, 2018 - Tartu, Estonia

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