

Haar Wavelet Method for Solving Stiff Differential Equations

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Abstract. Application of the Haar wavelet approach for solving stiff differential equations is discussed. Solution of singular perturbation problems is also considered. Efficiency of the recommended method is demonstrated by means of four numerical examples, mostly taken from well-known textbooks.

Key words: stiff differential equations, haar wavelets, numerical solution, nonlinear problems, singular perturbation problems.

1 Introduction

In the past two decades, stiff differential equations have been studied extensively and various methods for their solutions have been proposed. While the intuitive meaning of stiffness is clear to all specialists, a unique mathematical definition is lacking.

As a rule stiffness occurs in differential equations where there are two or more different time scales of the independent variables on which the dependent variables are changing. In the case of linear problems the stiffness is caused by eigenvalues of big negative values. To measure the degree of stiffness the following stiffness ratio $SR = \max |\lambda| / \min |\lambda|$ is introduced: when $SR < 20$ the problem is not stiff, up to $SR \approx 1000$ the problem is classified as stiff, and when $SR \geq 100000$, the problem is very stiff. Stiff ODEs are called extremely stable if there is at least one eigenvalue with a large negative real part [2].

In the case of nonlinear problems the problem is more complicated since stiffness is a global problem and cannot be reduced to the solution structure in the neighbourhood of single points.

If we apply for the solution of stiff differential equations explicit solvers (e.g. MATLAB solvers ODE45, ODE23s) the solution may be oscillating or diverge, where oscillations of the solution are a result of the failure of the solution algorithm. For avoiding these oscillations implicit solvers or explicit

solvers with a very small step size must be used. All this complicates the numerical work, besides small steps introduce too many round-off errors and cause numerical instabilities. Depending on the nature of stiffness different methods of solution should be used.

A detailed analysis of the methods for solving stiff ordinary differential equations can be found in the well known text-book [7] by Hairer and Wanner (here also an extensive list of publications is added). We would like to turn attention also to the paper [4] by Enright *et al.* in which different numerical methods have been tested for solving 25 systems of stiff equations. From the recent literature we refer here [1, 5, 6, 8, 12, 16, 18].

Of interest are also singular perturbation problems, where the differential equations contain a small parameter ε so that for $\varepsilon \rightarrow 0$ one of the equations loses its highest derivative [7]. In the limit case $\varepsilon = 0$ we get the so called differential algebraic system, which has been also investigated in many papers (see, e.g. [3, 5, 21]).

We are interested in solving stiff equations by the wavelet method. Hsiao [9, 10, 11] has proposed for solving linear and nonlinear differential equations the so called single-term Haar wavelet method (in fact it is a method of piecewise constant approximation). The method is very simple and effective, but its exactness in the region of rapid changes may turn out to be insufficient (see, Fig. 1 in [9]).

From recently published papers let us cite [19, 20]. In the these papers the numerical solutions of stiff differential equations, especially the Robertson problem, by using wavelet-based methods are investigated.

The main goal of the paper is to develop further Hsiao's results and demonstrate that the multi-term Haar wavelets are a powerful tool for solving stiff problems. The paper is organized as follows. In Section 2 the Haar wavelet method is shortly referred. In Section 3 this method is applied for solving linear stiff ODEs. Nonlinear problems are discussed in Section 4. Section 5 is devoted to solution of the Robertson's problem. Singular perturbation problems are considered in Section 6. In Section 7 some recommendations for further research are given.

2 Haar Wavelet Method

Let us consider the interval $x \in [A, B]$, where A and B are given constants. We shall define the quantity $M = 2^J$, where J is the maximal level of resolution. The interval $[A, B]$ is divided into $2M$ subintervals of equal length; the length of each subinterval is $\Delta x = (B - A)/(2M)$. Next two parameters are introduced: the dilatation parameter $j = 0, 1, \dots, J$ and the translation parameter $k = 0, 1, \dots, m - 1$ (here the notation $m = 2^j$ is introduced). The wavelet number i is identified as $i = m + k + 1$.

The i -th Haar wavelet is defined as

$$h_i(x) = \begin{cases} 1 & \text{for } x \in [\xi_1(i), \xi_2(i)], \\ -1 & \text{for } x \in [\xi_2(i), \xi_3(i)], \\ 0 & \text{elsewhere,} \end{cases} \quad (2.1)$$

where

$$\begin{aligned}\xi_1(i) &= A + 2k\mu\Delta x, & \xi_2(i) &= A + (2k+1)\mu\Delta x, \\ \xi_3(i) &= A + 2(k+1)\mu\Delta x, & \mu &= M/m.\end{aligned}$$

The case $i = 1$ corresponds to the scaling function: $h_1(x) = 1$ for $x \in [A, B]$ and $h_1(x) = 0$ elsewhere.

If we want to solve a n -th order ODE we need the integrals

$$p_{\nu,i}(x) = \underbrace{\int_A^x \int_A^x \dots \int_A^x}_{\nu\text{-times}} h_i(t) dt^\nu = \frac{1}{(\nu-1)!} \int_A^x (x-t)^{\nu-1} h_i(t) dt,$$

$$\nu = 1, 2, \dots, n, \quad i = 1, 2, \dots, 2M.$$

The case $\nu = 0$ corresponds to function $h_i(t)$. Taking into account (2.1), these integrals can be calculated analytically; by doing it we obtain

$$p_{\alpha,i}(x) = \begin{cases} 0 & \text{for } x < \xi_1(i), \\ \frac{1}{\alpha!} [x - \xi_1(i)]^\alpha & \text{for } x \in [\xi_1(i), \xi_2(i)], \\ \frac{1}{\alpha!} \{ [x - \xi_1(i)]^\alpha - 2[x - \xi_2(i)]^\alpha \} & \text{for } x \in [\xi_2(i), \xi_3(i)], \\ \frac{1}{\alpha!} \{ [x - \xi_1(i)]^\alpha - 2[x - \xi_2(i)]^\alpha + [x - \xi_3(i)]^\alpha \} & \text{for } x > \xi_3(i). \end{cases} \quad (2.2)$$

These formulas hold for $i > 1$. In the case $i = 1$ we have $\xi_1 = A$, $\xi_2 = \xi_3 = B$ and

$$p_{\alpha,1}(x) = \frac{1}{\alpha!} (x - A)^\alpha.$$

In the present paper the collocation method for solving the ODEs is applied. The collocation points are $x_l = 0.5[\tilde{x}_{l-1} + \tilde{x}_l]$, $l = 1, 2, \dots, 2M$: the symbol $\tilde{x}_l$ denotes the l -th grid point $\tilde{x}_l = A + l\Delta x$, $l = 1, 2, \dots, 2M$. Eqs. (2.1), (2.2) are discretized by replacing $x \rightarrow x_l$. It is convenient to introduce the Haar matrices $H(i, l) = h_i(x_l)$, $P_\nu(i, l) = p_{\nu,i}(x_l)$. In the following sections computer simulations are carried out with the aid of the *Matlab* programs for which the matrix representation is effective.

3 Linear Problems

If we have to solve a n -th order linear ODE we seek the solution in the form

$$y^{(n)}(x_l) = \sum_{i=1}^{2M} a_i h_i(x_l) = aH. \quad (3.1)$$

Lower order derivatives (and the function $y(x)$) are obtained through integrations of (3.1). All these ingredients are incorporated into (3.1) which is discretized by the collocation method. In this way we get a system of equations from which the wavelet coefficients a_i are calculated.

We investigate the accuracy of such algorithms. First, let us consider the case where exact solution $y_{ex}(x)$ of the problem is known. We define in the collocation points the error function $v(l) = y(x_l) - y_{ex}(x_l)$. Next the norm

$$\|v\|_p = \left(\sum_{l=1}^{2M} |v(l)|^p \right)^{1/p}$$

is introduced. The following two error estimates are applied (J denotes the resolution level):

- (i) local estimate $\delta_J = \left\| \frac{v}{y_{ex}} \right\|_{\infty} = \max_l \left| \frac{y(x_l)}{y_{ex}(x_l)} - 1 \right|$,
- (ii) global estimate $\sigma_J = \|v\|_2 / (2M)$.

Error estimates in the case when the exact solution is not known is considered further on in Sections 4–6.

Example 1. Press *et al.* [17] investigated the system

$$\begin{cases} y' = 998y + 1998z, & x \in [0, 1], \\ z' = -999y - 1999z, & y(0) = 1, \quad z(0) = 0. \end{cases} \quad (3.2)$$

We reduce (3.2) to the second order ODE

$$\begin{aligned} y'' + 1001y' + 1000y &= 0, & x \in [0, 1], \\ y(0) &= 1, & y'(0) = 0. \end{aligned} \quad (3.3)$$

The exact solution of (3.3) is

$$y_{ex} = \frac{1}{999} (1000e^{-x} - e^{-1000x}). \quad (3.4)$$

The eigenvalues of (3.3) are $\lambda_1 = -1$, $\lambda_2 = -1000$; since $|\lambda_2| \gg |\lambda_1|$ the system (3.2) is stiff. It is demonstrated in [17] that explicit methods of integration bring to oscillations and the solution is unstable.

Now let us solve (3.3) by the Haar wavelet method. The solution is sought in the form

$$\begin{aligned} y'' &= aH, \\ y' &= aP_1 + y'_0 E, \\ y &= aP_1 + y'_0 x + y_0 E. \end{aligned} \quad (3.5)$$

Here $E = (1, 1, \dots, 1)$ and all quantities are calculated in the collocation points. Substituting (3.5) into (3.3) we get the matrix equation

$$a(H + 1001P_1 + 1000P_2) = -1000E$$

from which the wavelet coefficients a_i are calculated. The solution of the problem $y = y(x)$ is found from (3.5)₃.

Computer simulations gave the following error estimates

$$\begin{aligned}\delta_3 &= 4.5e-4, & \sigma_3 &= 2.8e-5, \\ \delta_4 &= 1.7e-4, & \sigma_4 &= 5.3e-6, \\ \delta_5 &= 4.1e-5, & \sigma_5 &= 6.5e-7.\end{aligned}$$

It follows from here that high precision of the results is guaranteed already for a small number of collocation points (e.g. if $J = 3$ we have only 16 points).

Example 2. Mahmood *et al.* [16] solved by the Adomian decomposition method the linear initial value problem

$$\begin{cases} y_1' = -20y_1 - 0.25y_2 - 19.75y_3, & y_1(0) = 1 \\ y_2' = 20y_1 - 20.25y_2 + 0.25y_3, & y_2(0) = 0 \\ y_3' = 20y_1 - 19.75y_2 - 0.25y_3, & y_3(0) = -1 \end{cases} \quad (3.6)$$

which has an exact solution

$$\begin{aligned}y_{1ex} &= \frac{1}{2} [e^{-0.5t} + e^{-20t}(\cos 20t + \sin 20t)], \\ y_{2ex} &= \frac{1}{2} [e^{-0.5t} - e^{-20t}(\cos 20t - \sin 20t)], \\ y_{3ex} &= -\frac{1}{2} [e^{-0.5t} + e^{-20t}(\cos 20t - \sin 20t)].\end{aligned} \quad (3.7)$$

The solution (3.7) is plotted in Fig. 1.

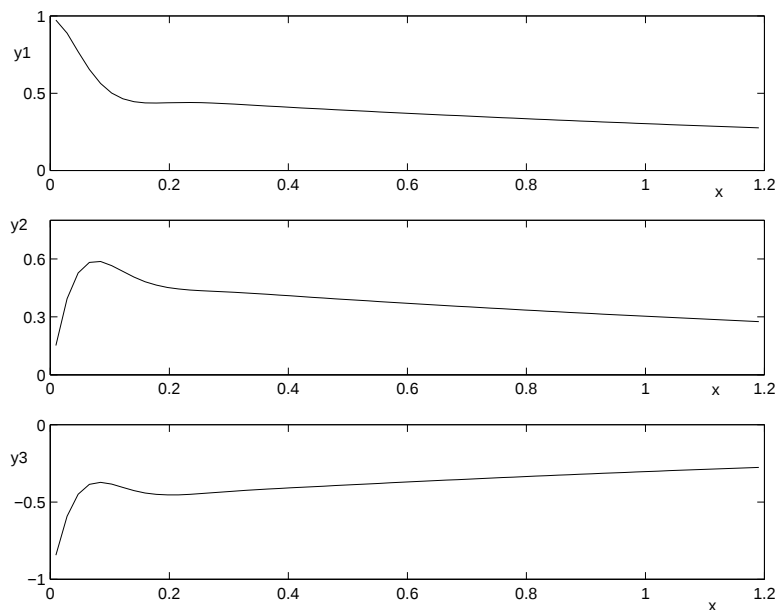


Figure 1. Solution of the system (3.6).

This problem was as well solved by Guzel and Bayram [6] who applied the power series method. Hojjati *et al.* [8] solved (3.4) with the aid of a predictor-corrector method, based on backward differentiation.

The wavelet solution is sought in the form

$$\begin{aligned}y_1' &= aH, & y_1 &= aP_1 + E \\y_2' &= bH, & y_2 &= bP_1 \\y_3' &= cH, & y_3 &= cP_1 - E.\end{aligned}\tag{3.8}$$

Substituting (3.8) into the system (3.6) we obtain the matrix system

$$\begin{aligned}a(H + 20P_1) + 0.25bP_1 + 19.75cP_1 &= -0.25E, \\-20aP_1 + b(H + 20.25P_1) - 0.25cP_1 &= 19.75E, \\-20aP_1 + 19.75bP_1 + c(H + 0.25P_1) &= 20.25E.\end{aligned}\tag{3.9}$$

Solving (3.9) we calculate the vectors a, b, c , the solution of the problem $y_1(x)$, $y_2(x)$, $y_3(x)$ is obtained from (3.8).

Computations for $J = 5$ (with 64 collocation points) were carried out and the results were compared with the exact solution (3.5). The following error estimates were obtained:

$$\begin{aligned}\delta_5(y_1) &= 2.0e - 5, & \sigma_5(y_1) &= 3.1e - 7, \\ \delta_5(y_2) &= 4.4e - 5, & \sigma_5(y_2) &= 6.9e - 7, \\ \delta_5(y_3) &= 1.6e - 3, & \sigma_5(y_3) &= 2.5e - 5.\end{aligned}$$

4 Nonlinear Problems

Consider the equation

$$y^{(n)} = f(x, y, y', \dots, y^{(n-1)}), \quad x \in [A, B]\tag{4.1}$$

with initial conditions $y^{(i)}(A) = y_0^i$, $i = 0, 1, \dots, 2M - 1$. Here f is a nonlinear function.

The wavelet solution of the problem is sought in the form (3.1). Calculating again lower order derivatives by integration, replacing these results into (4.1) and satisfying these equations in the collocation points we get a nonlinear system

$$F(x_l, a_1, \dots, a_{2M}) = 0, \quad l = 1, 2, \dots, 2M,\tag{4.2}$$

from which the wavelet coefficients a_i are calculated.

The system (4.2) is solved by the Newton method. If some initial guess for these coefficients $\tilde{a}_i$ is known, then the corrected values of the wavelet coefficients are computed as

$$a = \tilde{a} + \lambda \Delta a,$$

here $\Delta a = -F/\frac{\partial F}{\partial a}$ and λ is a coefficient, which is selected to guarantee the decrease of $\|F\|_\infty$ (in the case of the exact solution we have $\|F(a)\|_\infty = 0$). This Newton step is repeated until $\|F\|_\infty$ is sufficiently close to zero.

It is well known that the Newton method converges only when the initial guess is sufficiently good. To find such an initial estimate may be rather difficult. We have applied the following approach. We begin with a small number of collocation points (1, 2 or 4 points) for which the solution of (4.2) is more simple. Let us assume that we have found the solutions $a_i^{(0)}$ for some level of resolution J . The wavelet coefficients at the next level $J + 1$ are estimated as

$$\tilde{a}_i = \begin{cases} a_i^{(0)}, & \text{for } i = 1, 2, \dots, 2M, \\ 0, & \text{for } i = 2M + 1, 2M + 2, \dots, 4M. \end{cases}$$

Such an estimation is motivated by the fact that higher order coefficients of the sequence a_i are usually small (compare e.g. Fig. 2). These estimates are corrected by the Newton method and after that we can step over to the next level of resolution. Other possibilities for the initial guess of the wavelet coefficients are considered in Sections 5–6. For further details about this approach consult [13].

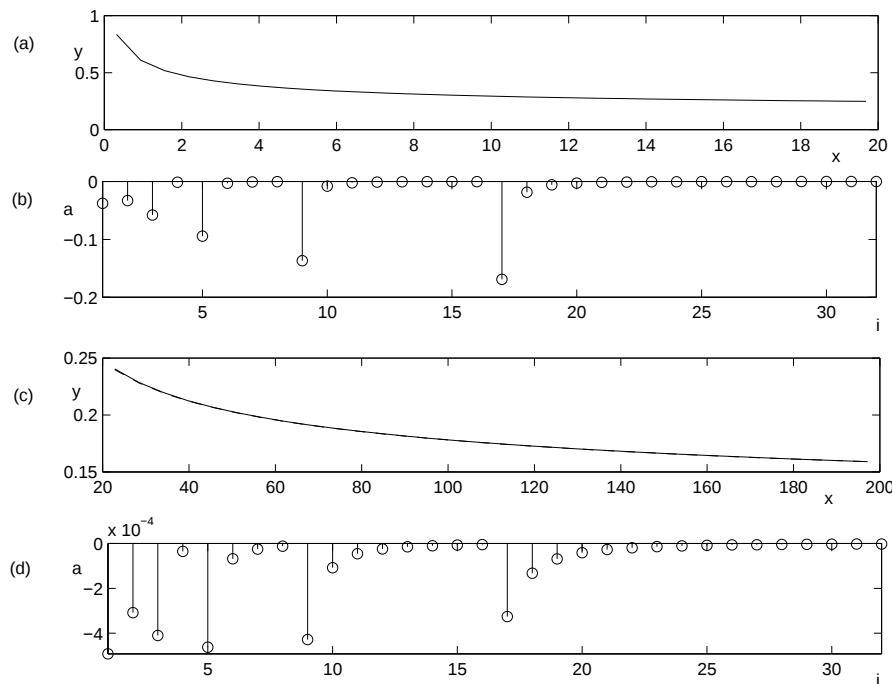


Figure 2. Solution of (4.3) for $J = 4$ (a) first phase $x \in [0, 20]$; (b) wavelet coefficients $a(i)$ for this phase; (c) second phase $x \in [20, 200]$; (d) wavelet coefficients for the second phase.

Example 3. Consider the initial value problem

$$y' = -\frac{y^2}{1+x}, \quad y(0) = 1, \quad x \in [0, 200], \quad (4.3)$$

which has the exact solution $y = 1/(1 + \ln(1 + x))$. This problem is classified by Hsiao [10] as a problem of "strong instability". Hsiao's solution is based on the single-term Haar wavelet method.

Let us solve (4.3) with the aid of multi-term Haar wavelets. At the beginning, the solution $y(x)$ rapidly decreases and then calms down. Therefore it is reasonable to solve the problem in two stages: first consider the phase of rapid decrease $x \in [0, \delta]$ and after that find solution for $x \in [\delta, 200]$. In the following calculations the value $\delta = 20$ was chosen. In both regions we take

$$y' = aH, \quad y = aP_1 + y_{in}E. \quad (4.4)$$

The function (4.2) and its gradient get the form

$$F(x) = aH + y/(E + x), \quad \frac{\partial F}{\partial a} = H + 2y \cdot P_1/(E + x).$$

Here $x = (x_1, x_2, \dots, x_{2M})$ and the point denotes element-by-element multiplication:

$$(y \cdot^2)_l = y(l)^2, \quad (y \cdot P_1)_{i,l} = y(l)P_1(i, l).$$

In the first phase $x \in [0, 20]$ we have $y_{in} = 1$. For the initial solution we take $a^0 \equiv 0$. This value is corrected by the Newton method. By increasing step by step the resolution level J the solution is made more and more exact. Computer simulation for $J = 5$ gave the error estimates $\delta_5 = 6.9e - 3$, $\sigma_5 = 6.0e - 5$. The value $y(20) = 0.2475$ was calculated, it is taken for the initial value y_{in} for the second stage $x \in [20, 200]$. It is assumed $a^0 \equiv 0$ and this value is again corrected by the Newton method. Here the number of collocation points may be smaller since already $\delta_3 = 8.5e - 3$, $\sigma_3 = 1.6e - 4$. The wavelet solution $y = y(x)$ and the wavelet coefficients a_i , $i = 1, 2, \dots, 2M$ are plotted in Fig. 2. It follows from this figure that the number of significant wavelet coefficients is rather small.

5 Robertson's Problem

In 1966 Robertson investigated a chemical system containing fast and slow motions at the same time. By modelling this system he got the following mathematical model

$$\begin{cases} y_1' = -0.04y_1 + 10^4 y_2 y_3, & y_1(0) = 1, \\ y_2' = 0.04y_1 - 10^4 y_2 y_3 - 3 \cdot 10^7 y_2^2, & y_2(0) = 0, \\ y_3' = 3 \cdot 10^7 y_2^2, & y_3(0) = 0. \end{cases} \quad (5.1)$$

The eigenvalues of (5.1) are given in [1]: $\lambda_1 = 0$, $\lambda_2 = -0.405$, $\lambda_3 = -2189.6$ and the third one produces the stiffness.

The system (5.1) has been investigated in several papers; it is often taken as a benchmark for numerical methods. A detailed analysis about it can be found in the text-book [7]. It is shown that explicit methods give oscillating solutions (see Fig. 1.3 in [7]). Stability was guaranteed in the case of implicit methods

(in [7] the Runge-Kutta codes DOPRI 5 and RADAU 5 were applied). In [11] the system (5.1) is integrated by the single term Haar wavelet method and in [14] the Adomian decomposition method was used. Solution obtained by the ODE45 code is plotted in Fig. 3. Small oscillations of the curve y_2 indicate instability of the solution.

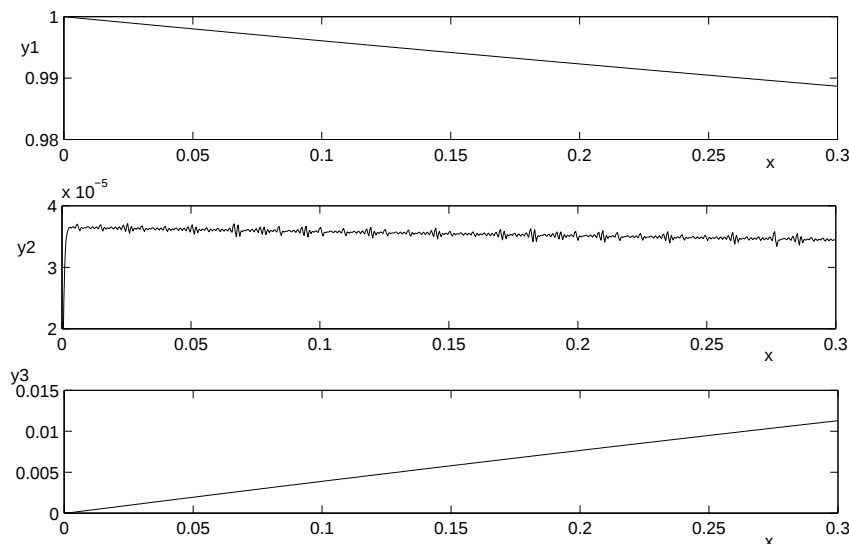


Figure 3. Solution of the Robertson's problem (4.4) obtained with the aid of the program ODE45.

Qualitative analysis of (5.1) indicates that the component y_2 rapidly reaches its maximal value for which $y_2' = 0$ and $3 \cdot 10^7 y_2^2 \approx 0.04$, consequently $\max(y_2) \approx 3.65e - 5$; after that the function y_2 decreases very slowly. Due to this fact it is again expedient to solve (5.1) separately for the regions $x \in [0, \delta]$ and $x \in [\delta, x_{max}]$. The wavelet solution is sought in the form

$$\begin{aligned} y_1' &= aH, & y_2' &= bH, & y_3' &= cH, \\ y_1 &= y_1(0)E + aP_1, & y_2 &= y_2(0)E + bP_1, & y_3 &= y_3(0)E + cP_1. \end{aligned} \quad (5.2)$$

Next the vectors F_1, F_2, F_3 are introduced

$$\begin{aligned} F_1 &= y_1' + 0.04y_1 - 10^4 y_2 \cdot y_3, \\ F_2 &= y_2' - 0.04y_1 + 10^4 y_2 \cdot y_3 + 3 \cdot 10^7 y_2^2, & F_3 &= y_3' - 3 \cdot 10^9 y_2^2. \end{aligned}$$

On account of (5.2) these vectors are functions of the wavelet coefficients a, b, c . We have to find such values a, b, c for which $F_1 = F_2 = F_3 = 0$. This can be realized according to the following procedure.

- (i) We assume that some initial estimates $\tilde{y}_1, \tilde{y}_2, \tilde{y}_3$ are known. From (5.2) we calculate the estimates $\tilde{a}, \tilde{b}, \tilde{c}$ and also $\tilde{y}_1', \tilde{y}_2', \tilde{y}_3'$.

- (ii) Evaluate F_2 .
- (iii) Since a rapid growth takes place for the function y_2 , we shall vary only the coefficients b (a and c are fixed) and calculate the gradient

$$\frac{\partial F_2}{\partial b} = H + 10^4 y_3 \cdot P_1 + 6 \cdot 10^7 y_2 \cdot P_1. \quad (5.3)$$

- (iv) Using the Newton method next approximation for b is calculated:

$$b = \tilde{b} - \lambda F_2 \left(\frac{\partial F_2}{\partial b} \right)^{-1}.$$

For the coefficient λ the value which minimizes $\|F_2\|$ is taken.

- (v) Calculate from (5.1)₃ and (5.2) the corrected values for c, y_3, y'_3 .
- (vi) Calculate from (5.1)₁ and (5.2) the corrected values for a, y_1, y'_1 .
- (vii) Evaluate $\|F_1\|_\infty, \|F_2\|_\infty, \|F_3\|_\infty$; if these norms are not sufficiently near to zero repeat the steps (ii)–(vii).

Computer simulations were carried out for $\delta = 0.005$, $x_{max} = 0.3$. In the first phase $x \in [0, \delta]$ the argument x is very small and we can develop the function y_1, y_2, y_3 into power series. Taking into account only the first terms of this series we find

$$\tilde{y}_1 = E - 0.04x, \quad \tilde{y}_2 = 0.04E - 1.6 \cdot 10^4 x.^3, \quad \tilde{y}_3 = 1.6 \cdot 10^4 x.^3.$$

According to (5.2) we obtain

$$\tilde{a} = -0.04E/H, \quad \tilde{b} = -4.8 \cdot 10^4 x.^2/H, \quad \tilde{c} = -\tilde{b}.$$

In order to correct these values the Newton iteration steps are done. Restricting ourselves to two iterations we get for $J = 5$:

$$y_1(\delta) = 0.9998, \quad y_2(\delta) = 3.648e - 5, \quad y_3(\delta) = 1.626e - 5.$$

These values were taken as initial values for the second phase $x \in [0.005, 0.3]$. In this phase the functions y_1, y_3 decrease very slowly, therefore we take

$$\tilde{y}_1 = y_1(\delta)E, \quad \tilde{y}_2 = y_2(\delta)E.$$

Calculating $\tilde{y}_3$ from (5.1)₃ we find

$$\tilde{y}_3 = y_3(\delta)E + 3 \cdot 10^7 y_2(\delta)^2 x.$$

Correcting once again the values by the Newton method we find

$$y_1(0.3) = 0.9888, \quad y_2(0.3) = 3.446e - 5, \quad y_3(0.3) = 1.138e - 4.$$

Carrying out the calculations by the ODE45 code we get

$$y_1(0.3) = 0.9987, \quad y_2(0.3) = 3.444e - 5, \quad y_3(0.3) = 1.129e - 4.$$

These results are in accordance with our results.

It is interesting to note that in the second phase only the first wavelet coefficients are significant. To demonstrate this, the first three wavelet coefficients are written down:

$$\begin{aligned} a : & -0.0376, \quad -0.0011, \quad -0.0006, \quad \dots \\ 10^5 b : & -0.6851, \quad -0.0183, \quad -0.0092, \quad \dots \\ c : & 0.0377, \quad 0.0011, \quad 0.0006, \quad \dots \end{aligned}$$

All the other coefficients are considerably smaller.

6 Singular Perturbation Problems

A singular perturbation problem of Index 1 is defined as [7]

$$\begin{cases} y' = f(y, z), & y(x_0) = y_0, \\ \varepsilon z' = g(y, z), & z(x_0) = z_0. \end{cases} \quad (6.1)$$

If ε is small, then the second equation of (5.3) is stiff. In the case $\varepsilon = 0$ we get a differential-algebraic equation, when a differential equation is combined with an algebraic equation. Eqs. (6.1) for $\varepsilon = 0$ are called also the *reduced system*. It is usually much easier to analyze the reduced system than the primary system.

Consider the phase space (y, z) . The slope of the phase curves is

$$\frac{dz}{dy} = \frac{g(y, z)}{\varepsilon f(y, z)}.$$

We assume that the initial point (y_0, z_0) is not placed on the curve $g(y, z) = 0$ and $\varepsilon \ll 1$; in this case the phase curve is very steep. We are mainly interested in the case where the phase point (y, z) moves towards the curve $g(y, z) = 0$ and reaches it at some instant $x = \delta$. Now $z'(\delta) = 0$ and the following motion proceeds along the curve $g(y, z) = 0$. Numerical difficulties can appear in the first phase of motion $x \leq \delta$, since for $\varepsilon \ll 1$ the second equation of (6.1) is stiff.

Example 4. Consider the van der Pol equation

$$\varepsilon z'' + (z^2 - 1)z' + z = 0,$$

which can be rewritten in the form [7]

$$\begin{cases} y' = -z, & y(x_0) = y_0, \\ \varepsilon z' = y - z^3/3 + z, & z(x_0) = z_0. \end{cases} \quad (6.2)$$

Boundary condition for z' can be calculated from (6.2)₂:

$$\varepsilon z'(x_0) = \left(y_0 - z_0^3/3 + z_0\right). \quad (6.3)$$

The reduced problem $\varepsilon = 0$ can be easily solved:

$$\ln |z| - 0.5z^2 = x + c.$$

Computer simulations were carried out for $\varepsilon = 0.001$, $y_0 = 0.5$, $z_0 = 1$. According to (6.3) we have $z'_0 = 7000/6$. Since $z' > 0$ the phase point (y, z) moves toward the phase curve

$$y = z^3/3 - z \quad (6.4)$$

and reaches it at some instant $x = \delta$.

Numerical results obtained by the MATLAB program ODE23s and with the aid of the piecewise constant approximation method (PCA) are plotted in Fig. 4. It follows from this figure that at this number of calculation points the PCA solution is unstable.

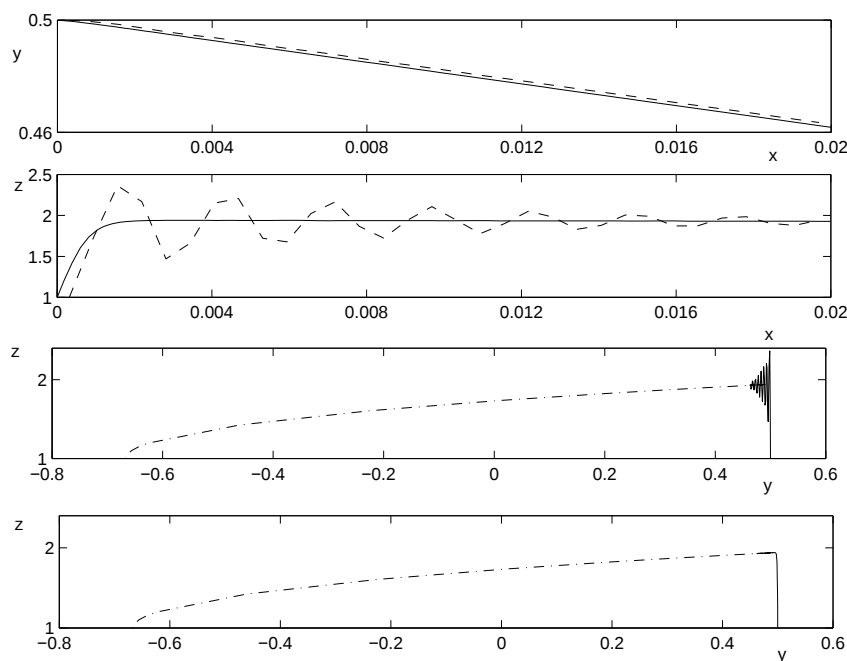


Figure 4. Solution of the van der Pol equation (6.1); (a)–(b) functions $y(x), z(x)$, dashed lines denote the PCA solutions; (c) phase curve according to the PCA solutions; (d) phase curve for the wavelet solution.

Now let us present the Haar wavelet solution. Again it is suitable to consider separately two phases of motion: (i) the phase of rapid changes $x \in [0, \delta]$ and (ii) motion along the curve (6.4) for $x \in [\delta, x_{fin}]$. In the following calculations it was taken $\delta = 0.003$, $x_{fin} = 0.02$.

In the first phase the solution is sought in the form

$$z'' = aH, \quad z' = aP_1 + z'_0 E, \quad z = aP_2 + z'_0 x + z_0 E.$$

Next the $2M$ -dimensional vector F is defined:

$$F(a) = \varepsilon z'' + (z.^2 - E).z' + z;$$

its gradient is given by

$$\frac{\partial F}{\partial a} = \varepsilon H + (z.^2 - E).P_1 + (E + 2z.z').P_2.$$

The corrected estimate for wavelet coefficients is given as

$$a = \tilde{a} - \lambda F \left(\frac{\partial F}{\partial a} \right)^{-1}.$$

Computer simulations were carried out for $J = 4$, for the initial approximation the value $\tilde{a} \equiv 0$ was taken and three Newton iteration steps were made. For the final result we got $\|F\|_\infty = 0.0026$, $y(\delta) = 0.4938$, $z(\delta) = 1.940$, $z'(\delta) = 0.9224$.

In the second phase $x \in [\delta, x_{fin}]$ we assume

$$z' = bH, \quad z = bP_1 + z(\delta)E.$$

Now the function

$$G(b) = (z.^2 - E).z' + z \quad (6.5)$$

is to be minimized. The gradient of (6.5) is given by

$$\frac{\partial G}{\partial b} = (z.^2 - E).H + (E + 2z.z').P_1.$$

Again we assume $\tilde{b} \equiv 0$; after two Newton steps we find $\|G\|_\infty = 8.5e - 5$, $y(0.02) = 0.4629$, $z(0.02) = 1.928$. This is in a good accordance with the ODE23s solution which gives $y(0.02) = 0.4618$, $z(0.02) = 1.928$.

As to the wavelet coefficients then the first coefficient is $b_1 = -0.7058$; all other coefficients are very small (less than 0.27% from b_1). This is due to the quasi-linearity of the function $z(x)$.

7 Conclusions

The examples of this paper demonstrate that in solving stiff systems the Haar wavelet method can successfully compete with the other efficient methods. The main benefits of the Haar approach are simplicity (already a small number of grid points guarantees the necessary accuracy) and universality (the same approach is applicable for a wide class of differential and integral equations). The subprograms, developed for calculation of integrals $p_{\alpha,i}(x)$ can be used without changes for solving different problems.

In the case of stiff problems the Haar wavelet method seems to be more stable than many other methods. There are some complementary possibilities to raise the stability of the solution. First, we could use the wavelets with a variable step size (see, [15], especially the example in Section 5). If we develop

into the Haar series the highest derivative $y^{(n)}$, then it is not continuous: this fact may also cause some instabilities. In such cases it is advisable to put into series the higher derivative $y^{(n+1)}$; this idea is implemented in [14].

In this paper only ordinary differential equations were considered, but the same approach is applicable also for stiff partial differential equations. We recommend to consult the paper [15] in which the Burgers equation was solved.

The Haar wavelet transform and PCA (single term Haar wavelet series) both consist of piecewise constant functions and therefore have the same convergence rate $O(M^{-2})$. Nevertheless the multi-Haar wavelet method has by our mind some advantages:

- (i) Very high accuracy fast transformation exists and therefore a possibility of implementation of fast algorithms compared with other methods;
- (ii) The simplicity and small computation costs, resulting from the sparsity of the transform matrices and small number of significant wavelet coefficients;
- (iii) The method is also very convenient for solving boundary value problems, since the boundary conditions are taken into account automatically.

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