

Canonical variational completion of 4D Gauss-Bonnet gravity

arXiv:2009.05459 [gr-qc]

Manuel Hohmann

Laboratory of Theoretical Physics - Institute of Physics - University of Tartu
Center of Excellence "The Dark Side of the Universe"



Virtual Conference of the Polish Society on Relativity - 26. 9. 2020

- 1 4D Gauss-Bonnet gravity
- 2 Canonical variational completion
- 3 Results
- 4 Conclusion

- 1 4D Gauss-Bonnet gravity
- 2 Canonical variational completion
- 3 Results
- 4 Conclusion

Gauss-Bonnet gravity

- Action for Gauss-Bonnet gravity in D dimensions:

$$S = \int d^D x \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \Lambda_0 + \alpha \mathcal{G} \right] + S_m. \quad (1)$$

Gauss-Bonnet gravity

- Action for Gauss-Bonnet gravity in D dimensions:

$$S = \int d^D x \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \Lambda_0 + \alpha \mathcal{G} \right] + S_m. \quad (1)$$

- Gauss-Bonnet scalar:

$$\mathcal{G} = 6R^{\mu\nu}{}_{[\mu\nu} R^{\rho\sigma}{}_{\rho\sigma]} = R^2 - 4R_{\mu\nu} R^{\mu\nu} + R_{\mu\nu\rho\sigma} R^{\mu\nu\rho\sigma}. \quad (2)$$

Gauss-Bonnet gravity

- Action for Gauss-Bonnet gravity in D dimensions:

$$S = \int d^D x \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \Lambda_0 + \alpha \mathcal{G} \right] + S_m. \quad (1)$$

- Gauss-Bonnet scalar:

$$\mathcal{G} = 6R^{\mu\nu}{}_{[\mu\nu} R^{\rho\sigma}{}_{\rho\sigma]} = R^2 - 4R_{\mu\nu} R^{\mu\nu} + R_{\mu\nu\rho\sigma} R^{\mu\nu\rho\sigma}. \quad (2)$$

⇒ Resulting field equations:

$$E_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} - 2\alpha \mathcal{G}_{\mu\nu} = T_{\mu\nu}. \quad (3)$$

Gauss-Bonnet gravity

- Action for Gauss-Bonnet gravity in D dimensions:

$$S = \int d^D x \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \Lambda_0 + \alpha \mathcal{G} \right] + S_m. \quad (1)$$

- Gauss-Bonnet scalar:

$$\mathcal{G} = 6R^{\mu\nu}{}_{[\mu\nu} R^{\rho\sigma}{}_{\rho\sigma]} = R^2 - 4R_{\mu\nu} R^{\mu\nu} + R_{\mu\nu\rho\sigma} R^{\mu\nu\rho\sigma}. \quad (2)$$

⇒ Resulting field equations:

$$E_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} - 2\alpha \mathcal{G}_{\mu\nu} = T_{\mu\nu}. \quad (3)$$

- Contribution from Gauss-Bonnet term:

$$\begin{aligned} \mathcal{G}_{\mu\nu} &= 15g_{\mu[\nu} R^{\rho\sigma}{}_{\rho\sigma} R^{\omega\tau}{}_{\omega\tau]} \\ &= \frac{1}{2} \mathcal{G} g_{\mu\nu} - 2R_{\mu\lambda\rho\sigma} R_{\nu}{}^{\lambda\rho\sigma} + 4R_{\mu\rho\nu\sigma} R^{\rho\sigma} + 4R_{\mu\rho} R_{\nu}{}^{\rho} - 2RR_{\mu\nu}. \end{aligned} \quad (4)$$

Gauss-Bonnet gravity

- Action for Gauss-Bonnet gravity in D dimensions:

$$S = \int d^D x \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \Lambda_0 + \alpha \mathcal{G} \right] + S_m. \quad (1)$$

- Gauss-Bonnet scalar:

$$\mathcal{G} = 6R^{\mu\nu}{}_{[\mu\nu} R^{\rho\sigma}{}_{\rho\sigma]} = R^2 - 4R_{\mu\nu} R^{\mu\nu} + R_{\mu\nu\rho\sigma} R^{\mu\nu\rho\sigma}. \quad (2)$$

⇒ Resulting field equations:

$$E_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} - 2\alpha \mathcal{G}_{\mu\nu} = T_{\mu\nu}. \quad (3)$$

- Contribution from Gauss-Bonnet term:

$$\begin{aligned} \mathcal{G}_{\mu\nu} &= 15g_{\mu[\nu} R^{\rho\sigma}{}_{\rho\sigma} R^{\omega\tau}{}_{\omega\tau]} \\ &= \frac{1}{2} \mathcal{G} g_{\mu\nu} - 2R_{\mu\lambda\rho\sigma} R_{\nu}{}^{\lambda\rho\sigma} + 4R_{\mu\rho\nu\sigma} R^{\rho\sigma} + 4R_{\mu\rho} R_{\nu}{}^{\rho} - 2RR_{\mu\nu}. \end{aligned} \quad (4)$$

⚡ Gauss-Bonnet contribution vanishes identically in $D = 4$.

4D Gauss-Bonnet gravity?

- Renormalization of Gauss-Bonnet term [Glavan, Lin '20]:

$$S = \int d^D x \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \Lambda_0 + \frac{\alpha}{D-4} \mathcal{G} \right] + S_m. \quad (5)$$

4D Gauss-Bonnet gravity?

- Renormalization of Gauss-Bonnet term [Glavan, Lin '20]:

$$S = \int d^D x \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \Lambda_0 + \frac{\alpha}{D-4} \mathcal{G} \right] + S_m. \quad (5)$$

⇒ Resulting field equations:

$$E_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} - \frac{2\alpha}{D-4} \mathcal{G}_{\mu\nu} = T_{\mu\nu}. \quad (6)$$

4D Gauss-Bonnet gravity?

- Renormalization of Gauss-Bonnet term [Glavan, Lin '20]:

$$S = \int d^D x \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \Lambda_0 + \frac{\alpha}{D-4} \mathcal{G} \right] + S_m. \quad (5)$$

⇒ Resulting field equations:

$$E_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} - \frac{2\alpha}{D-4} \mathcal{G}_{\mu\nu} = T_{\mu\nu}. \quad (6)$$

- Strategy for obtaining solutions in $D = 4$:
 - Solve field equations in arbitrary dimension D .
 - Consider limit $D \rightarrow 4$ on the solution.

4D Gauss-Bonnet gravity?

- Renormalization of Gauss-Bonnet term [Glavan, Lin '20]:

$$S = \int d^D x \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \Lambda_0 + \frac{\alpha}{D-4} \mathcal{G} \right] + S_m. \quad (5)$$

⇒ Resulting field equations:

$$E_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} - \frac{2\alpha}{D-4} \mathcal{G}_{\mu\nu} = T_{\mu\nu}. \quad (6)$$

- Strategy for obtaining solutions in $D = 4$:
 - Solve field equations in arbitrary dimension D .
 - Consider limit $D \rightarrow 4$ on the solution.
- Highly symmetric solutions have been obtained:
 - Homogeneity and isotropy ⇒ cosmological spacetimes.
 - Spherical symmetry ⇒ black holes, dust collapse etc.

4D Gauss-Bonnet gravity?

- Renormalization of Gauss-Bonnet term [Glavan, Lin '20]:

$$S = \int d^D x \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \Lambda_0 + \frac{\alpha}{D-4} \mathcal{G} \right] + S_m. \quad (5)$$

⇒ Resulting field equations:

$$E_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} - \frac{2\alpha}{D-4} \mathcal{G}_{\mu\nu} = T_{\mu\nu}. \quad (6)$$

- Strategy for obtaining solutions in $D = 4$:
 - Solve field equations in arbitrary dimension D .
 - Consider limit $D \rightarrow 4$ on the solution.
- Highly symmetric solutions have been obtained:
 - Homogeneity and isotropy ⇒ cosmological spacetimes.
 - Spherical symmetry ⇒ black holes, dust collapse etc.

⇒ **Non-vanishing contribution to solutions even in $D = 4$.**

- Structure of the Gauss-Bonnet contribution:

$$-\mathcal{G}_{\mu\nu} = (D - 4)A_{\mu\nu} + W_{\mu\nu}. \quad (7)$$

- Structure of the Gauss-Bonnet contribution:

$$-\mathcal{G}_{\mu\nu} = (D - 4)A_{\mu\nu} + W_{\mu\nu}. \quad (7)$$

- ✓ First term $\propto (D - 4)$ can be renormalized:

$$A_{\mu\nu} = \frac{D - 3}{(D - 2)^2} \left[\frac{2D}{D - 1} RR_{\mu\nu} - 4 \frac{D - 2}{D - 3} R^{\rho\lambda} C_{\mu\rho\nu\lambda} - 4R_{\mu}{}^{\rho} R_{\nu\rho} + 2R_{\rho\lambda} R^{\rho\lambda} g_{\mu\nu} - \frac{1}{2} \frac{D + 2}{D - 1} R^2 g_{\mu\nu} \right]. \quad (8)$$

- Structure of the Gauss-Bonnet contribution:

$$-\mathcal{G}_{\mu\nu} = (D - 4)A_{\mu\nu} + W_{\mu\nu}. \quad (7)$$

- ✓ First term $\propto (D - 4)$ can be renormalized:

$$A_{\mu\nu} = \frac{D - 3}{(D - 2)^2} \left[\frac{2D}{D - 1} RR_{\mu\nu} - 4 \frac{D - 2}{D - 3} R^{\rho\lambda} C_{\mu\rho\nu\lambda} - 4R_{\mu}{}^{\rho} R_{\nu\rho} + 2R_{\rho\lambda} R^{\rho\lambda} g_{\mu\nu} - \frac{1}{2} \frac{D + 2}{D - 1} R^2 g_{\mu\nu} \right]. \quad (8)$$

- ⚡ Second term not renormalizable (no factor $D - 4$ extractable):

$$W_{\mu\nu} = 2C_{\mu}{}^{\rho\lambda\sigma} C_{\nu\rho\lambda\sigma} - \frac{1}{2} C_{\tau\rho\lambda\sigma} C^{\tau\rho\lambda\sigma} g_{\mu\nu}. \quad (9)$$

- Structure of the Gauss-Bonnet contribution:

$$-\mathcal{G}_{\mu\nu} = (D - 4)A_{\mu\nu} + W_{\mu\nu}. \quad (7)$$

- ✓ First term $\propto (D - 4)$ can be renormalized:

$$A_{\mu\nu} = \frac{D - 3}{(D - 2)^2} \left[\frac{2D}{D - 1} RR_{\mu\nu} - 4 \frac{D - 2}{D - 3} R^{\rho\lambda} C_{\mu\rho\nu\lambda} - 4R_{\mu}{}^{\rho} R_{\nu\rho} + 2R_{\rho\lambda} R^{\rho\lambda} g_{\mu\nu} - \frac{1}{2} \frac{D + 2}{D - 1} R^2 g_{\mu\nu} \right]. \quad (8)$$

- ⚡ Second term not renormalizable (no factor $D - 4$ extractable):

$$W_{\mu\nu} = 2C_{\mu}{}^{\rho\lambda\sigma} C_{\nu\rho\lambda\sigma} - \frac{1}{2} C_{\tau\rho\lambda\sigma} C^{\tau\rho\lambda\sigma} g_{\mu\nu}. \quad (9)$$

- ? Truncate field equations and omit $W_{\mu\nu}$:

$$\dot{E}_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} + 2\alpha A_{\mu\nu} = T_{\mu\nu}. \quad (10)$$

- Structure of the Gauss-Bonnet contribution:

$$-\mathcal{G}_{\mu\nu} = (D - 4)A_{\mu\nu} + W_{\mu\nu}. \quad (7)$$

- ✓ First term $\propto (D - 4)$ can be renormalized:

$$A_{\mu\nu} = \frac{D - 3}{(D - 2)^2} \left[\frac{2D}{D - 1} RR_{\mu\nu} - 4 \frac{D - 2}{D - 3} R^{\rho\lambda} C_{\mu\rho\nu\lambda} - 4R_{\mu}{}^{\rho} R_{\nu\rho} + 2R_{\rho\lambda} R^{\rho\lambda} g_{\mu\nu} - \frac{1}{2} \frac{D + 2}{D - 1} R^2 g_{\mu\nu} \right]. \quad (8)$$

- ⚡ Second term not renormalizable (no factor $D - 4$ extractable):

$$W_{\mu\nu} = 2C_{\mu}{}^{\rho\lambda\sigma} C_{\nu\rho\lambda\sigma} - \frac{1}{2} C_{\tau\rho\lambda\sigma} C^{\tau\rho\lambda\sigma} g_{\mu\nu}. \quad (9)$$

- ? Truncate field equations and omit $W_{\mu\nu}$:

$$\dot{E}_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} + 2\alpha A_{\mu\nu} = T_{\mu\nu}. \quad (10)$$

- ⚡ Truncated equations are not variational for *any* D : $\nabla^{\mu} A_{\mu\nu} \neq 0$.

- Structure of the Gauss-Bonnet contribution:

$$-\mathcal{G}_{\mu\nu} = (D - 4)A_{\mu\nu} + W_{\mu\nu}. \quad (7)$$

- ✓ First term $\propto (D - 4)$ can be renormalized:

$$A_{\mu\nu} = \frac{D - 3}{(D - 2)^2} \left[\frac{2D}{D - 1} RR_{\mu\nu} - 4 \frac{D - 2}{D - 3} R^{\rho\lambda} C_{\mu\rho\nu\lambda} - 4R_{\mu}{}^{\rho} R_{\nu\rho} + 2R_{\rho\lambda} R^{\rho\lambda} g_{\mu\nu} - \frac{1}{2} \frac{D + 2}{D - 1} R^2 g_{\mu\nu} \right]. \quad (8)$$

- ⚡ Second term not renormalizable (no factor $D - 4$ extractable):

$$W_{\mu\nu} = 2C_{\mu}{}^{\rho\lambda\sigma} C_{\nu\rho\lambda\sigma} - \frac{1}{2} C_{\tau\rho\lambda\sigma} C^{\tau\rho\lambda\sigma} g_{\mu\nu}. \quad (9)$$

- ? Truncate field equations and omit $W_{\mu\nu}$:

$$\dot{E}_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} + 2\alpha A_{\mu\nu} = T_{\mu\nu}. \quad (10)$$

- ⚡ Truncated equations are not variational for *any* D : $\nabla^{\mu} A_{\mu\nu} \neq 0$.

- ? Try to find correction term to make them variational.

Outline

- 1 4D Gauss-Bonnet gravity
- 2 Canonical variational completion**
- 3 Results
- 4 Conclusion

Canonical variational completion

- Consider system of partial differential equations (PDEs):

$$\mathcal{E}_A(x^\mu, y^B, y^B_{\mu}, y^B_{\mu\nu}) = 0. \quad (11)$$

Canonical variational completion

- Consider system of partial differential equations (PDEs):

$$\mathcal{E}_A(x^\mu, y^B, y^B_{\mu}, y^B_{\mu\nu}) = 0. \quad (11)$$

- Variables: independent x^μ and dependent $y^B, y^B_{\mu} = \partial_{\mu}y^B \dots$

Canonical variational completion

- Consider system of partial differential equations (PDEs):

$$\mathcal{E}_A(x^\mu, y^B, y^B_{\mu}, y^B_{\mu\nu}) = 0. \quad (11)$$

- Variables: independent x^μ and dependent $y^B, y^B_{\mu} = \partial_{\mu}y^B \dots$
- Vainberg-Tonti Lagrangian density:

$$\mathcal{L}_{\mathcal{E}} = y^A \int_0^1 \mathcal{E}_A(x^\mu, ty^B, ty^B_{\mu}, ty^B_{\mu\nu}) dt. \quad (12)$$

Canonical variational completion

- Consider system of partial differential equations (PDEs):

$$\mathcal{E}_A(x^\mu, y^B, y^B_{\mu}, y^B_{\mu\nu}) = 0. \quad (11)$$

- Variables: independent x^μ and dependent $y^B, y^B_{\mu} = \partial_{\mu}y^B \dots$
- Vainberg-Tonti Lagrangian density:

$$\mathcal{L}_{\mathcal{E}} = y^A \int_0^1 \mathcal{E}_A(x^\mu, ty^B, ty^B_{\mu}, ty^B_{\mu\nu}) dt. \quad (12)$$

- Variationally completed PDE are Euler-Lagrange equations:

$$\tilde{\mathcal{E}}_A = \frac{\partial \mathcal{L}_{\mathcal{E}}}{\partial y^A} - d_{\mu} \frac{\partial \mathcal{L}_{\mathcal{E}}}{\partial y^A_{\mu}} + d_{\mu} d_{\nu} \frac{\partial \mathcal{L}_{\mathcal{E}}}{\partial y^A_{\mu\nu}}. \quad (13)$$

Canonical variational completion

- Consider system of partial differential equations (PDEs):

$$\mathcal{E}_A(x^\mu, y^B, y^B_{\mu}, y^B_{\mu\nu}) = 0. \quad (11)$$

- Variables: independent x^μ and dependent $y^B, y^B_{\mu} = \partial_{\mu}y^B \dots$
- Vainberg-Tonti Lagrangian density:

$$\mathcal{L}_{\mathcal{E}} = y^A \int_0^1 \mathcal{E}_A(x^\mu, ty^B, ty^B_{\mu}, ty^B_{\mu\nu}) dt. \quad (12)$$

- Variationally completed PDE are Euler-Lagrange equations:

$$\tilde{\mathcal{E}}_A = \frac{\partial \mathcal{L}_{\mathcal{E}}}{\partial y^A} - d_{\mu} \frac{\partial \mathcal{L}_{\mathcal{E}}}{\partial y^A_{\mu}} + d_{\mu} d_{\nu} \frac{\partial \mathcal{L}_{\mathcal{E}}}{\partial y^A_{\mu\nu}}. \quad (13)$$

⇒ Idea of *canonical variational completion* [Voicu, Krupka '15]:

- Original system \mathcal{E}_A is variational if and only if $\mathcal{E}_A = \tilde{\mathcal{E}}_A$.
- Otherwise, $H_A = \tilde{\mathcal{E}}_A - \mathcal{E}_A$ is canonical correction term.

Example: Einstein-Hilbert Lagrangian

1. Start from “incorrect Einstein equations” in $D = 4$

$$E_{\mu\nu} = -M_{\text{P}}^2 R_{\mu\nu} = T_{\mu\nu}. \quad (14)$$

Example: Einstein-Hilbert Lagrangian

1. Start from “incorrect Einstein equations” in $D = 4$

$$E_{\mu\nu} = -M_{\text{P}}^2 R_{\mu\nu} = T_{\mu\nu}. \quad (14)$$

2. Match factors and index position to definition of $T_{\mu\nu}$:

$$T^{\mu\nu} = \frac{2}{\sqrt{-g}} \frac{\delta \mathcal{L}_m}{\delta g_{\mu\nu}} \Rightarrow \mathcal{E}^{\mu\nu} = -\frac{1}{2} \sqrt{-g} E^{\mu\nu} = \frac{M_{\text{P}}^2}{2} \sqrt{-g} R^{\mu\nu}. \quad (15)$$

Example: Einstein-Hilbert Lagrangian

1. Start from “incorrect Einstein equations” in $D = 4$

$$E_{\mu\nu} = -M_{\text{P}}^2 R_{\mu\nu} = T_{\mu\nu}. \quad (14)$$

2. Match factors and index position to definition of $T_{\mu\nu}$:

$$T^{\mu\nu} = \frac{2}{\sqrt{-g}} \frac{\delta \mathcal{L}_m}{\delta g_{\mu\nu}} \Rightarrow \mathcal{E}^{\mu\nu} = -\frac{1}{2} \sqrt{-g} E^{\mu\nu} = \frac{M_{\text{P}}^2}{2} \sqrt{-g} R^{\mu\nu}. \quad (15)$$

3. Behavior of terms under scaling of variables $g_{\mu\nu} \rightarrow t g_{\mu\nu}$:

$$R_{\mu\nu} \rightarrow R_{\mu\nu}, \quad R^{\mu\nu} \rightarrow t^{-2} R^{\mu\nu}, \quad \sqrt{-g} \rightarrow t^2 \sqrt{-g}. \quad (16)$$

Example: Einstein-Hilbert Lagrangian

1. Start from “incorrect Einstein equations” in $D = 4$

$$E_{\mu\nu} = -M_{\text{P}}^2 R_{\mu\nu} = T_{\mu\nu}. \quad (14)$$

2. Match factors and index position to definition of $T_{\mu\nu}$:

$$T^{\mu\nu} = \frac{2}{\sqrt{-g}} \frac{\delta \mathcal{L}_m}{\delta g_{\mu\nu}} \Rightarrow \mathcal{E}^{\mu\nu} = -\frac{1}{2} \sqrt{-g} E^{\mu\nu} = \frac{M_{\text{P}}^2}{2} \sqrt{-g} R^{\mu\nu}. \quad (15)$$

3. Behavior of terms under scaling of variables $g_{\mu\nu} \rightarrow t g_{\mu\nu}$:

$$R_{\mu\nu} \rightarrow R_{\mu\nu}, \quad R^{\mu\nu} \rightarrow t^{-2} R^{\mu\nu}, \quad \sqrt{-g} \rightarrow t^2 \sqrt{-g}. \quad (16)$$

4. Vainberg-Tonti Lagrangian density:

$$\mathcal{L} = \int_0^1 t^0 \frac{M_{\text{P}}^2}{2} \sqrt{-g} R^{\mu\nu} g_{\mu\nu} dt = \frac{M_{\text{P}}^2}{2} \sqrt{-g} R. \quad (17)$$

Example: Einstein-Hilbert Lagrangian

1. Start from “incorrect Einstein equations” in $D = 4$

$$E_{\mu\nu} = -M_{\text{P}}^2 R_{\mu\nu} = T_{\mu\nu}. \quad (14)$$

2. Match factors and index position to definition of $T_{\mu\nu}$:

$$T^{\mu\nu} = \frac{2}{\sqrt{-g}} \frac{\delta \mathcal{L}_m}{\delta g_{\mu\nu}} \Rightarrow \mathcal{E}^{\mu\nu} = -\frac{1}{2} \sqrt{-g} E^{\mu\nu} = \frac{M_{\text{P}}^2}{2} \sqrt{-g} R^{\mu\nu}. \quad (15)$$

3. Behavior of terms under scaling of variables $g_{\mu\nu} \rightarrow t g_{\mu\nu}$:

$$R_{\mu\nu} \rightarrow R_{\mu\nu}, \quad R^{\mu\nu} \rightarrow t^{-2} R^{\mu\nu}, \quad \sqrt{-g} \rightarrow t^2 \sqrt{-g}. \quad (16)$$

4. Vainberg-Tonti Lagrangian density:

$$\mathcal{L} = \int_0^1 t^0 \frac{M_{\text{P}}^2}{2} \sqrt{-g} R^{\mu\nu} g_{\mu\nu} dt = \frac{M_{\text{P}}^2}{2} \sqrt{-g} R. \quad (17)$$

5. Completed field equations:

$$\tilde{E}_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} = T_{\mu\nu}. \quad (18)$$

Outline

- 1 4D Gauss-Bonnet gravity
- 2 Canonical variational completion
- 3 Results**
- 4 Conclusion

Vainberg-Tonti Lagrangian of Gauss-Bonnet gravity

- Start from field equations in arbitrary dimension:

$$E_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} + 2\alpha \left(A_{\mu\nu} + \frac{W_{\mu\nu}}{D-4} \right) = T_{\mu\nu}. \quad (19)$$

Vainberg-Tonti Lagrangian of Gauss-Bonnet gravity

- Start from field equations in arbitrary dimension:

$$E_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} + 2\alpha \left(A_{\mu\nu} + \frac{W_{\mu\nu}}{D-4} \right) = T_{\mu\nu}. \quad (19)$$

- Behavior of terms under scaling of variables $g_{\mu\nu} \rightarrow t g_{\mu\nu}$:

$$g_{\mu\nu} \rightarrow t g_{\mu\nu}, \quad G_{\mu\nu} \rightarrow G_{\mu\nu}, \quad \{A, W\}_{\mu\nu} \rightarrow t^{-1} \{A, W\}_{\mu\nu}. \quad (20)$$

Vainberg-Tonti Lagrangian of Gauss-Bonnet gravity

- Start from field equations in arbitrary dimension:

$$E_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} + 2\alpha \left(A_{\mu\nu} + \frac{W_{\mu\nu}}{D-4} \right) = T_{\mu\nu}. \quad (19)$$

- Behavior of terms under scaling of variables $g_{\mu\nu} \rightarrow t g_{\mu\nu}$:

$$g_{\mu\nu} \rightarrow t g_{\mu\nu}, \quad G_{\mu\nu} \rightarrow G_{\mu\nu}, \quad \{A, W\}_{\mu\nu} \rightarrow t^{-1} \{A, W\}_{\mu\nu}. \quad (20)$$

⇒ Vainberg-Tonti Lagrangian density [MH, Pfeifer, Voicu '20]:

$$\begin{aligned} \mathcal{L} &= -\frac{1}{2} g_{\mu\nu} \int_0^1 t^{D/2} \sqrt{-g} \\ &\quad \left[t^{-2} M_{\text{P}}^2 G^{\mu\nu} + t^{-1} \Lambda_0 g^{\mu\nu} + 2t^{-3} \alpha \left(A^{\mu\nu} + \frac{W^{\mu\nu}}{D-4} \right) \right] dt \quad (21) \\ &= \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \Lambda_0 - \frac{2\alpha}{D-4} \left(A^\mu{}_\mu + \frac{W^\mu{}_\mu}{D-4} \right) \right]. \end{aligned}$$

Vainberg-Tonti Lagrangian of Gauss-Bonnet gravity

- Start from field equations in arbitrary dimension:

$$E_{\mu\nu} = M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} + 2\alpha \left(A_{\mu\nu} + \frac{W_{\mu\nu}}{D-4} \right) = T_{\mu\nu}. \quad (19)$$

- Behavior of terms under scaling of variables $g_{\mu\nu} \rightarrow t g_{\mu\nu}$:

$$g_{\mu\nu} \rightarrow t g_{\mu\nu}, \quad G_{\mu\nu} \rightarrow G_{\mu\nu}, \quad \{A, W\}_{\mu\nu} \rightarrow t^{-1} \{A, W\}_{\mu\nu}. \quad (20)$$

⇒ Vainberg-Tonti Lagrangian density [MH, Pfeifer, Voicu '20]:

$$\begin{aligned} \mathcal{L} &= -\frac{1}{2} g_{\mu\nu} \int_0^1 t^{D/2} \sqrt{-g} \\ &\quad \left[t^{-2} M_{\text{P}}^2 G^{\mu\nu} + t^{-1} \Lambda_0 g^{\mu\nu} + 2t^{-3} \alpha \left(A^{\mu\nu} + \frac{W^{\mu\nu}}{D-4} \right) \right] dt \quad (21) \\ &= \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \Lambda_0 - \frac{2\alpha}{D-4} \left(A^\mu{}_\mu + \frac{W^\mu{}_\mu}{D-4} \right) \right]. \end{aligned}$$

⚡ Note term $\propto t^{D/2-3}$ yielding factor $(D/2 - 2)^{-1} = 2(D-4)^{-1}$.

Consistency check: full field equations $E_{\mu\nu}$ with $W_{\mu\nu}$

- Calculate traces:

$$A^\mu{}_\mu = \frac{D-3}{D-2} \left(2R_{\mu\nu}R^{\mu\nu} - \frac{DR^2}{2(D-1)} \right), \quad (22a)$$

$$W^\mu{}_\mu = (D-4) \left(\frac{2R_{\mu\nu}R^{\mu\nu}}{D-2} - \frac{R^2}{(D-1)(D-2)} - \frac{R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma}}{2} \right). \quad (22b)$$

Consistency check: full field equations $E_{\mu\nu}$ with $W_{\mu\nu}$

- Calculate traces:

$$A^\mu{}_\mu = \frac{D-3}{D-2} \left(2R_{\mu\nu}R^{\mu\nu} - \frac{DR^2}{2(D-1)} \right), \quad (22a)$$

$$W^\mu{}_\mu = (D-4) \left(\frac{2R_{\mu\nu}R^{\mu\nu}}{D-2} - \frac{R^2}{(D-1)(D-2)} - \frac{R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma}}{2} \right). \quad (22b)$$

⇒ Note that traces satisfy:

$$A^\mu{}_\mu + \frac{W^\mu{}_\mu}{D-4} = -\frac{1}{2}\mathcal{G}. \quad (23)$$

Consistency check: full field equations $E_{\mu\nu}$ with $W_{\mu\nu}$

- Calculate traces:

$$A^\mu{}_\mu = \frac{D-3}{D-2} \left(2R_{\mu\nu}R^{\mu\nu} - \frac{DR^2}{2(D-1)} \right), \quad (22a)$$

$$W^\mu{}_\mu = (D-4) \left(\frac{2R_{\mu\nu}R^{\mu\nu}}{D-2} - \frac{R^2}{(D-1)(D-2)} - \frac{R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma}}{2} \right). \quad (22b)$$

⇒ Note that traces satisfy:

$$A^\mu{}_\mu + \frac{W^\mu{}_\mu}{D-4} = -\frac{1}{2}\mathcal{G}. \quad (23)$$

✓ For $D > 4$, reproduce original Lagrangian density:

$$\mathcal{L} = \sqrt{-g} \left(\frac{M_{\text{P}}^2}{2} R - \Lambda_0 + \frac{\alpha}{D-4} \mathcal{G} \right). \quad (24)$$

Consistency check: full field equations $E_{\mu\nu}$ with $W_{\mu\nu}$

- Calculate traces:

$$A^\mu{}_\mu = \frac{D-3}{D-2} \left(2R_{\mu\nu}R^{\mu\nu} - \frac{DR^2}{2(D-1)} \right), \quad (22a)$$

$$W^\mu{}_\mu = (D-4) \left(\frac{2R_{\mu\nu}R^{\mu\nu}}{D-2} - \frac{R^2}{(D-1)(D-2)} - \frac{R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma}}{2} \right). \quad (22b)$$

⇒ Note that traces satisfy:

$$A^\mu{}_\mu + \frac{W^\mu{}_\mu}{D-4} = -\frac{1}{2}\mathcal{G}. \quad (23)$$

✓ For $D > 4$, reproduce original Lagrangian density:

$$\mathcal{L} = \sqrt{-g} \left(\frac{M_{\text{P}}^2}{2} R - \Lambda_0 + \frac{\alpha}{D-4} \mathcal{G} \right). \quad (24)$$

⚡ Action diverges for $D \rightarrow 4$.

Variational completion of truncated equations $\mathring{E}_{\mu\nu}$

- Consider Vainberg-Tonti Lagrangian of truncated field equations:

$$\mathring{\mathcal{L}} = \sqrt{-g} \left(\frac{M_{\text{P}}^2}{2} R - \Lambda_0 - \frac{2\alpha}{D-4} A^{\mu}{}_{\mu} \right). \quad (25)$$

Variational completion of truncated equations $\mathring{E}_{\mu\nu}$

- Consider Vainberg-Tonti Lagrangian of truncated field equations:

$$\mathring{\mathcal{L}} = \sqrt{-g} \left(\frac{M_{\text{P}}^2}{2} R - \Lambda_0 - \frac{2\alpha}{D-4} A^\mu{}_\mu \right). \quad (25)$$

⇒ For $D > 4$, obtain variational fourth-order equations:

$$\begin{aligned} \mathring{\mathring{E}}_{\mu\nu} = & M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} + \frac{4\alpha(D-3)}{(D-1)(D-2)(D-4)} \\ & \left[g_{\mu\nu} \left(\square R - DR^2/4 + (D-1)R^{\rho\sigma} R_{\rho\sigma} \right) - 2(D-1)\square R_{\mu\nu} \right. \\ & \left. + (D-2)\nabla_\mu \nabla_\nu R + DRR_{\mu\nu} - 4(D-1)R^{\rho\sigma} R_{\mu\rho\nu\sigma} \right] = T_{\mu\nu}. \quad (26) \end{aligned}$$

Variational completion of truncated equations $\mathring{E}_{\mu\nu}$

- Consider Vainberg-Tonti Lagrangian of truncated field equations:

$$\mathring{\mathcal{L}} = \sqrt{-g} \left(\frac{M_{\text{P}}^2}{2} R - \Lambda_0 - \frac{2\alpha}{D-4} A^\mu{}_\mu \right). \quad (25)$$

⇒ For $D > 4$, obtain variational fourth-order equations:

$$\begin{aligned} \mathring{\mathring{E}}_{\mu\nu} = & M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} + \frac{4\alpha(D-3)}{(D-1)(D-2)(D-4)} \\ & \left[g_{\mu\nu} \left(\square R - DR^2/4 + (D-1)R^{\rho\sigma} R_{\rho\sigma} \right) - 2(D-1)\square R_{\mu\nu} \right. \\ & \left. + (D-2)\nabla_\mu \nabla_\nu R + DRR_{\mu\nu} - 4(D-1)R^{\rho\sigma} R_{\mu\rho\nu\sigma} \right] = T_{\mu\nu}. \quad (26) \end{aligned}$$

⚡ For $D = 4$, Vainberg-Tonti Lagrangian diverges due to t^{-1} term:

$$\int_0^1 t^{D/2-3} dt \xrightarrow{D \rightarrow 4} \int_0^1 \frac{dt}{t} \rightarrow \infty. \quad (27)$$

Variational completion of truncated equations $\mathring{E}_{\mu\nu}$

- Consider Vainberg-Tonti Lagrangian of truncated field equations:

$$\mathring{\mathcal{L}} = \sqrt{-g} \left(\frac{M_{\text{P}}^2}{2} R - \Lambda_0 - \frac{2\alpha}{D-4} A^\mu{}_\mu \right). \quad (25)$$

⇒ For $D > 4$, obtain variational fourth-order equations:

$$\begin{aligned} \mathring{\mathring{E}}_{\mu\nu} = & M_{\text{P}}^2 G_{\mu\nu} + \Lambda_0 g_{\mu\nu} + \frac{4\alpha(D-3)}{(D-1)(D-2)(D-4)} \\ & \left[g_{\mu\nu} \left(\square R - DR^2/4 + (D-1)R^{\rho\sigma} R_{\rho\sigma} \right) - 2(D-1)\square R_{\mu\nu} \right. \\ & \left. + (D-2)\nabla_\mu \nabla_\nu R + DRR_{\mu\nu} - 4(D-1)R^{\rho\sigma} R_{\mu\rho\nu\sigma} \right] = T_{\mu\nu}. \quad (26) \end{aligned}$$

⚡ For $D = 4$, Vainberg-Tonti Lagrangian diverges due to t^{-1} term:

$$\int_0^1 t^{D/2-3} dt \xrightarrow{D \rightarrow 4} \int_0^1 \frac{dt}{t} \rightarrow \infty. \quad (27)$$

⚡ No variational completion for $D = 4$ Gauss-Bonnet gravity.

Outline

- 1 4D Gauss-Bonnet gravity
- 2 Canonical variational completion
- 3 Results
- 4 Conclusion**

- Objectives:
 - Considered Gauss-Bonnet field equations in arbitrary dimension D .
 - Considered both full and truncated (no $W_{\mu\nu}$ term) field equations.
 - Applied canonical variational completion procedure.

- Objectives:

- Considered Gauss-Bonnet field equations in arbitrary dimension D .
- Considered both full and truncated (no $W_{\mu\nu}$ term) field equations.
- Applied canonical variational completion procedure.

⇒ Results:

- ✓ Reproduced Gauss-Bonnet action from full equations in $D > 4$.
- ⇒ Truncated equations: 4'th order variational completion in $D > 4$.
- ⚡ Truncated equations are not variational in *any* dimension.
- ⚡ Neither equation can be variationally completed in $D = 4$.

Summary

- Objectives:

- Considered Gauss-Bonnet field equations in arbitrary dimension D .
- Considered both full and truncated (no $W_{\mu\nu}$ term) field equations.
- Applied canonical variational completion procedure.

⇒ Results:

- ✓ Reproduced Gauss-Bonnet action from full equations in $D > 4$.
- ⇒ Truncated equations: 4'th order variational completion in $D > 4$.
- ⚡ Truncated equations are not variational in *any* dimension.
- ⚡ Neither equation can be variationally completed in $D = 4$.

Further reading

MH, C. Pfeifer and N. Voicu,
“Canonical variational completion of 4D Gauss-Bonnet gravity”,
arXiv:2009.05459 [gr-qc].